

# Fine Tuning of Lithographic Masks through Thin Films of PS-*b*-PMMA with Different Molar Mass by Rapid Thermal Processing

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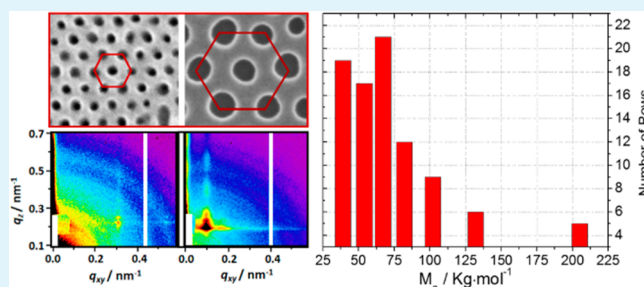
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**ABSTRACT:** The self-assembly of asymmetric polystyrene-*b*-poly(methyl methacrylate) (PS-*b*-PMMA) block copolymer based nanoporous thin films over a broad range of molar mass ( $M_n$ ) between 39 kg·mol<sup>-1</sup> and 205 kg·mol<sup>-1</sup> is obtained by means of a simple thermal treatment. In the case of standard thermal treatments, the self-assembly process of block copolymers is hindered at small  $M_n$  by thermodynamic limitations and by a large kinetic barrier at high  $M_n$ . We demonstrate that a fine tuning of the annealing parameters, performed by a Rapid Thermal Processing (RTP) machine, permits us to overcome those limitations. Cylindrical features are obtained by varying  $M_n$  and properly changing the corresponding annealing temperature, while keeping constant the annealing time (900 s), the film thickness (~30 nm), and the PS fraction (~0.7). The morphology, the characteristic dimensions (i.e., the pore diameter  $d$  and the pore-to-pore distance  $L_0$ ), and the order parameter (i.e., the lattice correlation length  $\xi$ ) of the samples are analyzed by scanning electron microscopy and grazing-incidence small-angle X-ray scattering, obtaining values of  $d$  ranging between 12 and 30 nm and  $L_0$  ranging between 24 and 73 nm. The dependence of  $L_0$  as a 0.67 power law of the number of segments places these systems inside the strong segregation limit regime. The experimental results evidence the capability to tailor the self-assembly processes of block copolymers over a wide range of molecular weights by a simple thermal process, fully compatible with the stringent constraints of lithographic applications and industrial manufacturing.

**KEYWORDS:** block copolymer, self-assembly, rapid thermal processing, GISAXS analysis, nanopatterning, molecular mass



## INTRODUCTION

Highly immiscible block copolymers (BCPs) have gained increasing interest over the last decade due to their ability to self-assemble into structures featuring periodicities at the nanometric scale. Due to this peculiar characteristic, BCPs appear to be suitable candidates for technologies involving surface patterning to be exploited in a wide range of emerging fields in nanotechnology, such as virus and water filtration,<sup>1–3</sup> organic optoelectronics,<sup>4</sup> anti-reflection coatings,<sup>5,6</sup> and nanolithography.<sup>7–9</sup> Current “top-down” photolithographic procedures are well established and highly optimized, supporting the progressive scaling-down of microelectronics over the last few decades. Nevertheless, inherent physical limitations related to optical diffraction in 193 nm lithography<sup>10</sup> forced micro-electronic industries to develop quite complex and expensive procedures, such as double patterning and immersion lithography,<sup>11</sup> in order to achieve the goals of the 32 and 22 nm nodes. A further extension of the photolithographic techniques beyond the 22 nm limit would require the

introduction of even more complex and expensive multiple exposure lithographic tools. In this context, non-optical “bottom up” lithographic approaches using self-assembling (SA) materials have been proposed as a complementary process to partially overcome the limitations of current photolithography.<sup>12,13</sup> In this regard, thin films of SA BCPs are of great interest for application in the semiconductor industry, owing to the possibility of integrating these cost-effective lithographic materials with the standard “top-down” processes.<sup>14</sup>

From a general point of view, the morphology of the BCP film, involving lamellae, gyroids, spheres, and cylinders,<sup>15</sup> can be controlled by varying the relative volume fraction  $f$  of the two blocks whereas, within each morphology, the structure periodicity can be tuned by varying the BCP molar mass ( $M_n$ ).

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In this context, several BCPs have been already investigated, such as polystyrene-*b*-polydimethylsiloxane (PS-*b*-PDMS), polystyrene-*b*-poly(methyl methacrylate) (PS-*b*-PMMA), polystyrene-*b*-poly(2-vinylpyridine) (PS-*b*-P2VP), or polystyrene-*b*-poly(ethylene oxide) (PS-*b*-PEO).

Among these, PS-*b*-PMMA represents a good candidate for integration with current “top down” approaches, because it holds most of the characteristics required by an ideal sub-lithographic patterning material.<sup>16</sup> First of all, PS-*b*-PMMA is soluble in common solvents such as toluene or tetrahydrofuran. Then, the selectivity of the PMMA with respect to acetic acid allows a selective removal of the PMMA phase, thus leading to the fabrication of a PS robust template for the pattern transfer to the substrate. Furthermore, due to the relatively weak interaction energy between the two blocks, both perpendicular and parallel orientation of the nanodomains are easily accessible by properly controlling the surface–polymer interaction<sup>17,18</sup> through the deposition of homopolymer<sup>19</sup> or random copolymer (RCP) thin films.<sup>20,21</sup> The main limitation of PS-*b*-PMMA is the minimum feature size obtainable with this BCP. The minimum achievable size of the periodicity of the polymeric film is dictated by the relatively low energy interaction between the two blocks, which prevents the phase separation in low  $M_n$  PS-*b*-PMMA BCPs. On the other side, the low mobility and slow ordering kinetics of the high  $M_n$  PS-*b*-PMMA BCPs control the maximum achievable size of the periodicity.

Extremely long annealing times (between  $10^2$  and  $10^3$  min) are required to obtain the vertical orientation, which represents the suited orientation for the preparation of a lithographic mask, when the self-assembling process is driven by pure thermal treatments.<sup>22,23</sup> For this reason, alternative or complementary techniques have been investigated including the use of external electric fields during the thermal treatment,<sup>24,25</sup> directional solvent vapor flow,<sup>26,27</sup> shear-induced alignment,<sup>28</sup> and solvent vapor annealing (SVA), eventually in combination with microwave radiation<sup>29</sup> or thermal annealing.<sup>30,31</sup> Nevertheless, the complexity of solvothermal setup and the dependence of the SA process on the experimental parameters of the SVA system (i.e., distance between solvent and sample, solvent vapor composition, absolute chamber pressure, temperature, humidity) result in poor reproducibility of the samples.<sup>32</sup>

Recently, a novel technology, rapid thermal processing (RTP), was employed to control the SA process of BCPs.<sup>33,34</sup> This technique involves the use of the infrared radiation of the RTP apparatus to govern the temperature of the polymeric films. Very high heating and cooling rates can be adopted, thus making it possible to attain high temperature short time thermal cycles. Well organized cylinder morphologies in films of 30 nm of an asymmetric  $67 \text{ kg mol}^{-1}$  PS-*b*-PMMA were obtained in less than 60 s.<sup>33</sup>

The present work extends the use of RTP to the processing of thin films of cylinder forming PS-*b*-PMMA samples featuring a wide range of molar mass, with  $M_n$  comprised between 39 to  $205 \text{ kg mol}^{-1}$ . In this context, it is pertinent to observe that whereas for symmetric BCPs the SA process dependence on the  $M_n$  has been intensively investigated,<sup>35–37</sup> only few works on this issue are found in the literature for asymmetric PS-*b*-PMMA thin films.<sup>24,38</sup>

## EXPERIMENTAL SECTION

**Surface Neutralization.** Oriented Si(100) substrates with a 100 nm thick thermal SiO<sub>2</sub> layer were used as a support for the BCP deposition. The oxidized substrates (about  $1 \text{ cm}^2$  surface) were cleaned with piranha solution (H<sub>2</sub>SO<sub>4</sub>/H<sub>2</sub>O<sub>2</sub> with 3/1 vol. ratio at 80 °C for 40 min) to eliminate any residual organic material and to increase the surface concentration of hydroxyl groups. The samples were then rinsed in H<sub>2</sub>O, dried under N<sub>2</sub> flow, cleaned with isopropanol in an ultrasonic bath, and finally dried under N<sub>2</sub> flow. After the cleaning process, a solution of a functional poly(styrene-*r*-methyl methacrylate) (P(*S-r*-MMA)) with styrene fraction 0.62 ( $M_n = 13.5 \text{ kg mol}^{-1}$  and polydispersity index PDI = 1.26, Polymer Source Inc.) (18.0 mg in 2.0 mL of toluene) was prepared in an ultrasonic bath and then spun on the substrates for 30 s at 3000 rpm. This particular RCP was chosen on the basis of previous results reported in the literature on the vertical orientation of the cylindrical nanodomains obtained on the same BCP.<sup>39,40</sup> In these conditions, a film thickness around 30 nm is obtained. Grafting was obtained by thermally treating the samples<sup>21</sup> in RTP under N<sub>2</sub> atmosphere. The RTP machine is a standard semiconductor manufacturing equipment that employs radiative energy sources (normally halogen lamps) to drive the desired temperature on the sample with extremely high heating rates. In the present experiments the grafting of the RCP was promoted with an annealing process at 250 °C for 600 s with a heating rate of 18 °C/s. The cooling velocity was controlled by changing the N<sub>2</sub> flow rate, resulting in a cooling rate of 10 °C/s. Consequently, the time required to bring the temperature of the sample below the  $T_g$  was about 20 s. After the removal of the non-grafted P(*S-r*-MMA), performed with an ultrasonic bath in toluene, the thickness of the grafted RCP layer was  $\sim 7$  nm.

**Block Copolymer Deposition and Self-Assembly.** Asymmetric PS-*b*-PMMA BCPs with different molar masses ( $M_n$ , styrene volume fraction  $f$ , and PDI values are listed in Table 1) were purchased from

**Table 1. Molar Mass ( $M_n$ ) of the BCP, Molar Mass ( $M_n$  Styrene) of the Polystyrene Block, Polystyrene Volume Fraction ( $f$ ), and Polydispersity Index (PDI) of All the BCPs**

sample	$M_n$ (kg/mol)	$M_n$ Styrene (kg/mol)	$f$	PDI
B29	29.4	21.5	0.75537	1.10
B39	39.0	26.8	0.71367	1.07
B54	53.8	37.0	0.71419	1.07
B67	67.1	46.1	0.71353	1.09
B82	82.0	57.0	0.72121	1.07
B102	101.5	68.0	0.69725	1.08
B132	132.0	96.0	0.75159	1.11
B205	205.0	140.0	0.70962	1.16

Polymer Source Inc. and used without further purification. The samples having different  $M_n$  have been marked with the letter B, standing for Block, and a number representing the corresponding value of molar mass (e.g. B102 corresponds to the BCP with  $M_n = 102 \text{ kg mol}^{-1}$ ). Thermogravimetric analyses (TGA) were performed using a Mettler-Toledo thermobalance mod. TGA/SDTA 851° LF1100 at a scanning rate of  $20 \text{ °C min}^{-1}$  from room temperature up to 600 °C under N<sub>2</sub> flow.

A solution of the various PS-*b*-PMMA (18.0 mg in 2.0 mL of toluene) was spun for 30 s at 3000 rpm on the grafted RCP. The self-organization of the PS-*b*-PMMA films was induced, by RTP under N<sub>2</sub> flow, heating the samples to the annealing temperature ( $T_A$ ) at  $18 \text{ °C s}^{-1}$  and maintaining the samples at  $T_A$  for 900 s.

The thickness of the polymeric films was measured by means of an M-200U spectroscopic ellipsometer (J. A. Wollam Co. Inc.) using light radiation from a Xe lamp impinging on the sample at a 70° incidence angle. Films of controlled thickness ( $\sim 30$  nm) were deposited for each PS-*b*-PMMA BCP on the properly neutralized substrates via spin-coating from dilute solutions of PS-*b*-PMMA in toluene (1% w/w) for 30 s at 3000 rpm.

**Morphological Analysis.** The morphology of the BCPs (in-plane order and orientation with respect to the substrate) was investigated by means of scanning electron microscopy (SEM, Zeiss Supra 40) and grazing-incidence small-angle X-ray scattering (GISAXS). GISAXS was performed at the ID01 beamline of the European Synchrotron Radiation Facility (ESRF, Grenoble, France), using an incident beam energy of 8.0 keV, corresponding to a wavelength  $\lambda = 0.155$  nm, and a two-dimensional MAXIPIX fast readout photon-counting pixel detector (four chips,  $516 \times 516$  pixels, pixel size  $55 \mu\text{m}$ ), positioned at the end of a vacuum tube. The beam diameter was  $\sim 100 \mu\text{m}$  and the sample-to-detector distance 0.830 m. Samples and detector were mounted on a high resolution  $\theta$ - $2\theta$  goniometer, allowing a fine control of the incidence angle.

In order to enhance the phase contrast in both SEM and GISAXS characterizations, a selective removal of the PMMA component was performed. The degradation of the PMMA and parallel cross-linking of the PS were achieved by exposure to UV radiation ( $5 \text{ mW cm}^{-2}$ ,  $\lambda = 253.7$  nm, 900 s). The PMMA material was subsequently removed through immersion in an acetic acid bath and subsequent rinsing in deionized  $\text{H}_2\text{O}$ . Finally, oxygen plasma treatments were performed to remove the RCP at the bottom of the pores.

## RESULTS AND DISCUSSION

In general, the deposition and processing of a specific BCP involve the optimization of several parameters including the film thickness as well as the annealing time and temperature. However, the constraints dictated by the successful exploitation of these materials in industrial applications impose certain limits which represent, to some extent, inescapable boundary conditions.

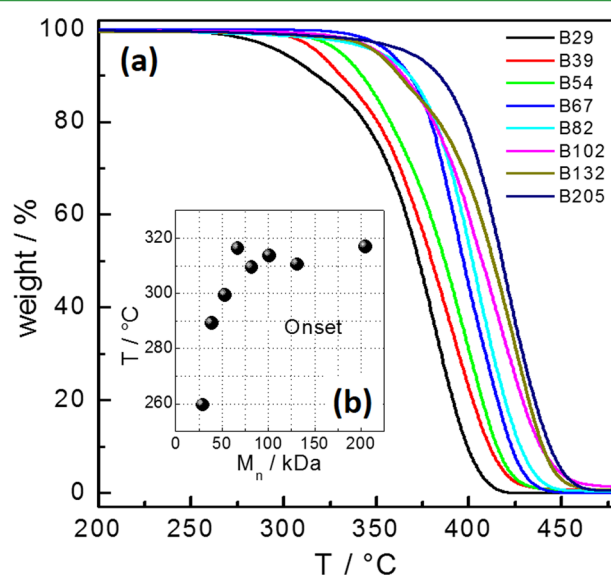
When the ordering process is carried out by conventional thermal treatments in an oven/furnace, the formation of uniform patterns usually requires combining proper optimization of the annealing conditions with fine tuning of the thickness of the BCP film.<sup>16</sup> The thickness to produce a uniform polymeric template over the sample surface decreases as the  $M_n$  of the BCP decreases. However, in view of the technological application of the PS-*b*-PMMA scaffold as sacrificial layer for pattern transfer, the thickness of the deposited film has to sustain additive or subtractive strategies for nanostructure fabrication on the underlying substrate. In this regard, the possibility of a nanoporous polymeric film of fixed thickness ( $\sim 30$  nm) with tunable characteristic lengths ( $d$ ,  $L_0$ ) would represent a step towards the simplification of the subsequent lithographic processes.

Even if the effectiveness of annealing times shorter than 60 s has already been demonstrated for relatively small  $M_n$  BCPs,<sup>33,34</sup> in this work an annealing time of 900 s was always adopted to afford all the BCPs, irrespective of their specific  $M_n$ , enough time to self-organize with a minimum level of order. Moreover, this annealing time allows maintaining the samples in an isothermal condition for a relatively long time at the different investigated temperatures. This prolonged isothermal condition allows minimizing the contribution of the transients on the overall thermal treatment and to effectively investigate the effect of  $T_A$  variations on the microphase separation and ordering process. Nevertheless, this annealing time is still much shorter than those adopted for other thermal treatments and definitely fulfills the industrial requirements for manufacturing processes. Consequently, the phase separation and lateral organization of the PS-*b*-PMMA BCPs with different  $M_n$  values (from  $29.4 \text{ kg}\cdot\text{mol}^{-1}$  to  $205 \text{ kg}\cdot\text{mol}^{-1}$ ) are promoted by changing the annealing temperature  $T_A$  in a RTP system. The surface neutralization, all the deposition (i.e., BCP film

thickness and type of solvent) and annealing parameters (i.e., annealing time, heating ramp, and  $\text{N}_2$  flow) were kept constant.

Once these boundary experimental conditions were defined, the optimization of the SA process mainly relies on the determination of the window of temperatures that can be operated to promote the phase separation and ordering of the BCP with the PMMA blocks arranged into hexagonally packed cylinders oriented perpendicularly to the substrate. The glass transition temperature ( $T_G$ ) sets the lower limit whereas the upper working temperature is given by the degradation temperature ( $T_{\text{DEG}}$ ) of the PS-*b*-PMMA. The value of the  $T_G$  in the bulk phase is known to be  $T_G \sim 110$  °C for the PS and  $T_G \sim 124$  °C for the PMMA, with a very limited variation ( $\Delta T_G \pm 2$  °C) depending on the  $M_n$ . On the other hand, the  $T_{\text{DEG}}$  of the BCPs are known to be significantly affected by the  $M_n$ . Accordingly, the thermal stability of each BCP was investigated through TGA analyses. Although the  $T_G$  and the  $T_{\text{DEG}}$  values in thin films are expected to differ from those measured in bulk, the measurement of the aforementioned limits in the bulk phase represents a useful guideline for the optimization of the  $T_A$ .

Figure 1a reports collectively the TGA curves, recorded under  $\text{N}_2$  flow at the heating rate of  $20$  °C $\cdot\text{min}^{-1}$ , for all the

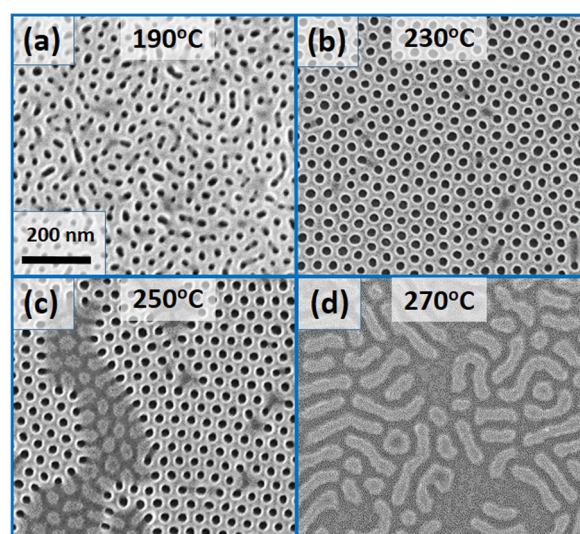


**Figure 1.** (a) TGA curves at  $20$  °C $\cdot\text{min}^{-1}$  heating rate under  $\text{N}_2$  flow for the PS-*b*-PMMA BCPs with different  $M_n$ . The inset (b) reports the temperatures corresponding to the onsets of the thermal degradation as a function of BCP  $M_n$ .

BCPs. In all cases, a sigmoidal shape of the loss curve is observed. As the  $M_n$  increases, the loss curves are regularly shifted toward higher temperature values. However, subtle differences can be noticed within the various samples, especially corresponding to the initial stages of the thermal degradation. For the lowest  $M_n$  sample (B29), the beginning of the loss curve occurs at the lowest temperature (about  $250$  °C, Figure 1b), and a single weight loss is observed extending over a wide temperature region. As the molar mass increases, the temperature corresponding to the initial degradation stage increases and the width of the weight loss profile slightly decreases. These differences in the thermal degradation profile probably account for the increasingly more relevant contribution of the end groups in initiating the thermal degradation, as

the average  $M_n$  decreases.<sup>41</sup> On the basis of the latter results we set the onset values  $T_{\text{DEG}}$  of the degradation curves as upper bounds of the working temperature window. The relevant  $T_{\text{DEG}}$  data as a function of  $M_n$  are reported in Figure 1b. The  $T_{\text{DEG}}$  values initially increase and subsequently level off at high  $M_n$ , further supporting the hypothesis of a significant role of the end groups in the thermal degradation process. It is worth noting that the  $T_{\text{DEG}}$  increase with the  $M_n$  leads to a progressive enlargement of the window of accessible temperatures. In turn, the  $T_{\text{DEG}}$  increase allows the high  $M_n$  BCPs to be processed at high temperatures, thus decreasing the melt viscosity and increasing the kinetics of the ordering process.

**Morphological Analysis.** To establish the optimum  $T_A$  range for the SA of each BCP, a systematic SEM analysis of the morphology of the annealed polymeric films was performed in a wide temperature range. Figure 2 shows a series of SEM



**Figure 2.** SEM images of sample B102 at four different annealing temperatures: (a) 190 °C, (b) 230 °C, (c) 250 °C, and (d) 270 °C.

images for sample B102 as a typical example. At 190 °C, the SA process appears in a preliminary stage (Figure 2a). The phase separation and the perpendicular orientation of the PMMA cylinders are achieved, but the hexagonal order of the cylinders is not completely developed.<sup>42</sup> In contrast, at 230 °C a well-defined hexagonal arrangement of the PMMA cylinders is obtained with a good lateral order (Figure 2b). On further increasing the  $T_A$  (250 °C), a number of defects become apparent (Figure 2c), and, finally, the hexagonal pattern

completely disappears for  $T_A$  higher than 270 °C (Figure 2d). It is worth noting that, under the adopted experimental conditions, the nanometric films do not show any sign of degradation. Therefore, the breakdown of the hexagonal arrangement at  $T_A \geq 250$  °C is likely to occur because the  $T_A$  approaches the order–disorder transition temperature ( $T_{\text{ODT}}$ ).

A similar analysis was performed on all the BCPs, and the relevant  $T_A$  values are collected in Table 2. The occurrence of nanodomains with hexagonally packed cylinders, normally oriented with respect to the substrate, was observed for all the samples with the only exception of B29. For the latter sample, phase separation and perpendicular orientation of the cylinders occur but the lateral order is very low, as shown in Figure 3a,b, where the SEM images of sample B39 and B29 are compared.

The above results can be discussed within the frame of the mean-field phase diagram reported in Figure 3c. In the PS-*b*-PMMA system, the formation of a regular pattern occurs when the product  $\chi N$  ( $\chi$  being the Flory–Huggins parameter and  $N$  the number of statistical segments per BCP chain) is higher than the minimum value that corresponds to the transition point from the ordered to the disordered state (ODT).<sup>8,43</sup> Pursuant to the phase diagram, no ordering can be achieved in lamellar forming symmetric BCPs with  $\chi N < 10.5$  and in cylinder forming asymmetric BCPs having  $\chi N < 15$ . Although for PS-*b*-PMMA  $\chi$  is only weakly temperature dependent, in the computation of the phase diagram, the  $\chi N$  values at the corresponding  $f$  were calculated for each  $M_n$  at the effective  $T_A$ . By superimposing the  $\chi N$  values obtained for the different BCPs on the phase diagram, we observe that, for all the samples except the B29 BCP, the  $\chi N$  values fall inside the area in which the cylindrical phase is expected to occur.

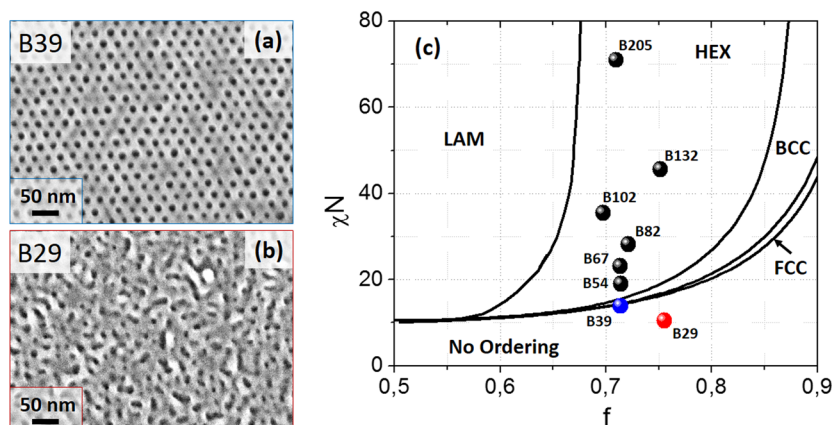
The  $\chi N$  value of the B29 ( $\chi N = \sim 10.5$ ) with  $f = 0.73$  lies well inside the zone of the phase diagram corresponding to the disordered phase. However, a clear evidence of phase separation with subsequent perpendicular orientation of the nanodomains was observed also for B29. This particular experimental behavior can be accounted for by considering that the phase diagram reported in Figure 3c was determined based on mean-field theory calculations for bulk materials, whereas the actual behavior of thin films should include the effect of the interfaces which in the present case seem to increase the propensity of the BCPs to segregate. Consequently, the reported phase diagram provides essentially a guideline for the interpretation of the results.

**Characteristic Dimensions.** Systematic SEM and GISAXS analyses were performed on all the samples exhibiting effective lateral organization to determine their morphological param-

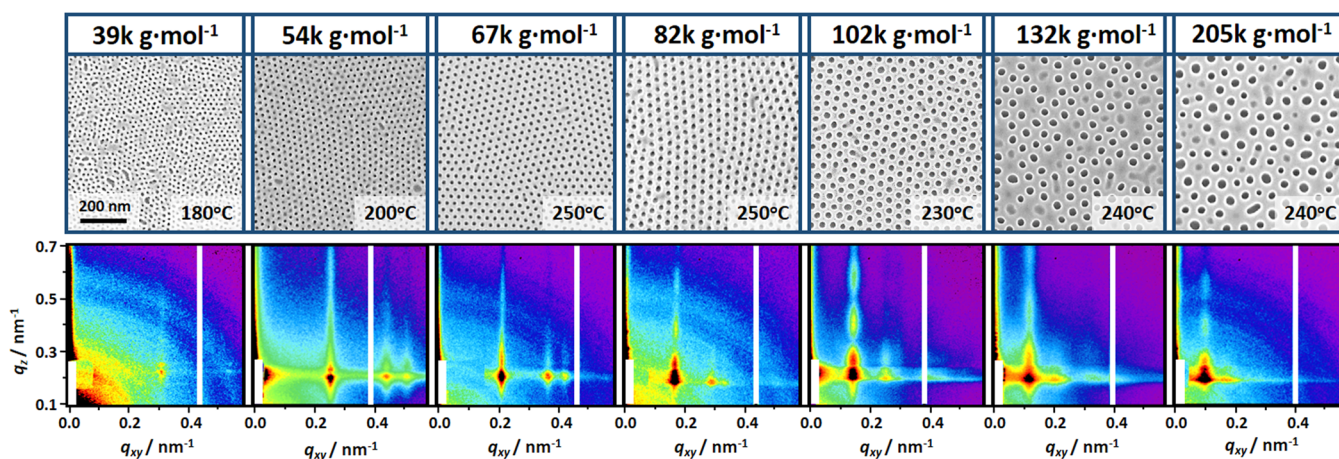
**Table 2. Preparation and Morphological Parameters of Optimized BCP Samples: Film Thickness  $h$  (measured by ellipsometry), Annealing Temperature  $T_A$ ,  $\chi N$  Values Calculated at  $T = T_A$ , Cylinder Diameter  $d$  (measured by SEM), Lattice Spacing Measured by SEM  $L_{\text{0SEM}}$  and by GISAXS  $L_{\text{0GISAXS}}$ , and Correlation Length  $\xi$  (measured by GISAXS)<sup>a</sup>**

sample	$h$ (nm)	$T_A$ (°C)	$\chi N$ (at $T_A$ )	$d$ (nm)	$L_{\text{0SEM}}$ (nm)	$L_{\text{0GISAXS}}$ (nm)	$h/L_0$	$\xi$ (nm)	$\xi/L_0$
B39	33	180	13.8	12.0 ± 2.0	24.0 ± 1.0	23.5 ± 0.2	1.37	447	19
B54	29	200	18.0	13.0 ± 1.0	28.8 ± 0.5	28.3 ± 0.3	1.00	477	17
B67	27	250	23.1	17.0 ± 1.0	35.0 ± 1.0	34.4 ± 0.4	0.77	710	21
B82	34	250	28.2	19.0 ± 2.0	42.9 ± 0.7	43.9 ± 0.7	0.79	535	12
B102	33	230	35.5	22.7 ± 1.5	47.0 ± 1.0	48.7 ± 0.9	0.70	434	9
B132	28	240	45.6	28.6 ± 1.6	59.0 ± 4.0	62.1 ± 1.4	0.47	354	6
B205	35	240	70.9	30.0 ± 4.0	73.0 ± 8.0	74.0 ± 2.0	0.48	357	5

<sup>a</sup>The value of  $L_0$  used to calculate the ratios  $h/L_0$  and  $\xi/L_0$  are the ones obtained by SEM.



**Figure 3.** SEM images representing the morphology of the B39 (blue circle) and that of the B29 (red circle) are shown in (a) and (b), respectively. (c) Phase diagram of the BCP representing the  $\chi N$  product (calculated for the specific  $T_A$ , reported in Table 2) as a function of the styrene volume fraction  $f$  (adapted from Figure 2 in ref 8 with additional data). The black lines represent the phase boundaries of the lamellar phase (LAM), hexagonally packed cylindrical phase (HEX), spherical phase (BCC and FCC), and disorder phase. The circles correspond to the BCPs listed in Table 1.



**Figure 4.** SEM images and corresponding GISAXS patterns for PS-*b*-PMMA thin films of increasing  $M_n$  after removal of PMMA domains. The corresponding annealing temperatures are also reported. All GISAXS patterns have been recorded at an incidence angle  $\alpha_i = 0.18^\circ$ , except the B39 where  $\alpha_i = 0.15^\circ$ . The white vertical line in the GISAXS patterns corresponds to the blind gap separating the active chips of the MAXIPIX detector.

ters. A collection of representative SEM images and corresponding GISAXS patterns is reported in Figure 4.

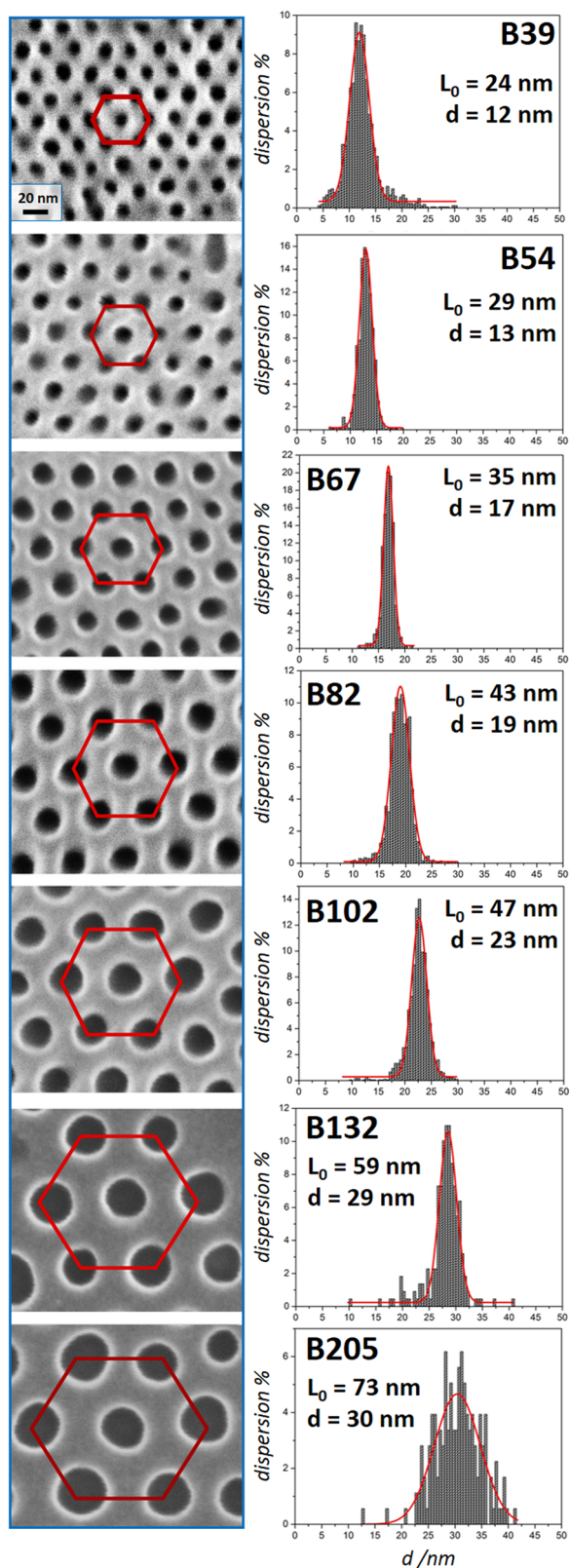
The SEM images show PMMA cylinders assembled to form a two-dimensional hexagonal lattice, with both the cylinder diameter  $d$  and the lattice constant  $L_0$  progressively increasing with  $M_n$ . A close view of the different hexagonal patterns is provided in the left column of Figure 5 for all the BCP thin films.

The statistical analysis of the SEM images allows determining the values of  $d$  and  $L_0$  with high accuracy. The histograms representing the statistical distribution of the PMMA cylinder diameters are reported in the right column of Figure 5: the dispersion in size of the PMMA cylinders is quite limited, apart from the B205 BCP. The broader dispersion of the cylinder size for this sample should be accounted for by considering both the higher polydispersity and the high molar mass which results in a very low chain mobility. The values of  $d$ , ranging from  $12 \pm 2$  nm to  $30 \pm 4$  nm, and  $L_0$ , from  $24 \pm 1$  nm to  $73 \pm 8$  nm, are reported in Table 2.

Contrary to what was observed in samples annealed in the furnace,<sup>44</sup> the RTP treatment affords perpendicular orientation of the PMMA cylinders irrespectively of the ratio  $h/L_0$ . Since  $h$

was nearly constant in the various samples, we obtained a different ratio  $h/L_0$  for each  $M_n$ , with values varying from  $h/L_0 \sim 4/3$  for sample B39 to  $h/L_0 \sim 1/2$  for sample B205, as reported in Table 2. The perpendicular orientation was observed even in samples having  $h/L_0 \sim 3/4$  (samples B82 and B102), in which the formation of half cylinders parallel to the substrate was observed to occur.<sup>44</sup>

The SEM morphological analysis was complemented by GISAXS investigation performed on all the samples. This technique allows exploring a much wider area than SEM ( $\sim 10^6$  times larger), providing additional information on the spatial extent of cylinder positional order and on its uniformity over the depth of the film.<sup>45,46</sup> Representative GISAXS patterns, taken above the critical angle ( $\alpha_c \approx 0.14^\circ$ ), are shown in Figure 4 below the corresponding SEM images. The quantities  $q_z$  and  $q_{xy}$  denote the components of the scattering vector normal and parallel to the plane of the sample, respectively. The presence of intensity rods, parallel to the  $q_z$  axis and set at discrete  $q_{xy}$  values, unequivocally indicates the in-plane periodicity of the PMMA domains and their orientation orthogonal to the plane of the film. The modulation of the rods along  $q_z$  present in all the patterns though more visible in high  $M_n$  samples, is related



**Figure 5.** SEM images of the BCP thin films (left column) and the pore diameter distribution (on the right). The  $L_0$  and  $d$  values indicated in the figure are rounded to the nearest integer. The exact values are reported in Table 2 with the corresponding error.

to the film thickness and provides values of  $h$  in agreement with those evaluated by ellipsometry and listed in Table 2.

Figure 6a shows the intensity profiles along  $q_{xy}$  obtained from horizontal cuts of the GISAXS patterns of Figure 4 at  $q_z = 0.22 \text{ nm}^{-1}$  (sample B39) and  $q_z = 0.24 \text{ nm}^{-1}$  (all other samples), i.e., on the maxima of scattered intensity. A series of peaks is visible, which can be indexed as the  $(h,k)$  reflections of a two-dimensional hexagonal lattice:

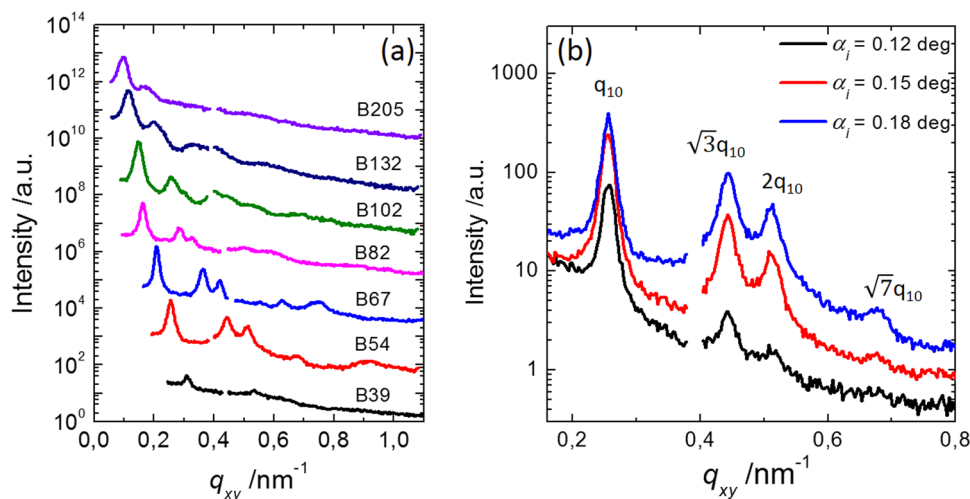
$$q_{h,k} = \frac{4\pi\sqrt{h^2 + hk + k^2}}{\sqrt{3}L_0}$$

Their shift towards lower  $q$  values is consistent with the anticipated increase of  $L_0$  for higher  $M_n$ . In particular, the values of lattice spacing calculated from the  $q$  vector of the first order (1,0) reflection,  $L_0 = 4\pi/\sqrt{3}q_{10}$ , are in very good agreement with those obtained by SEM (see Table 2). Further measurements were carried out both above and below the critical angle, to probe the structural properties of the surface layer at different depths over the whole film thickness: at  $\alpha_i = 0.12^\circ$  the radiation penetration depth was calculated to be  $\sim 10$  nm, i.e., one third of the film thickness  $h$ , whereas it was much larger than  $h$  for  $\alpha_i = 0.15^\circ$  and  $\alpha_i = 0.18^\circ$ . As shown in Figure 6b for the B54 sample, the comparison of the intensity profiles obtained from GISAXS patterns taken above and below the critical angle does not show significant differences in the peak position and width, thus suggesting that the film structure remains nearly uniform over the whole thickness of the sample. The presence of diffraction peaks at scattering vectors in the ratio  $1:\sqrt{3}:\sqrt{4}:\sqrt{7}$  confirms the hexagonal arrangement of the PMMA cylinders, oriented normal to the substrate.

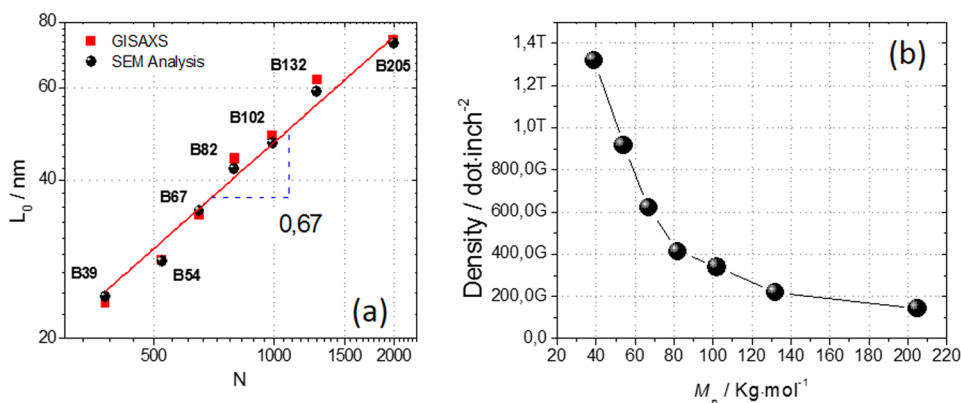
In Figure 7, the values of lattice spacing  $L_0$  obtained both by SEM and GISAXS analyses are reported as a function of the number of segments  $N$  in the copolymer chain. The two sets of values are in excellent agreement within the experimental error and  $L_0$  exhibits a clear exponential dependence on  $N$ , with  $L_0 \propto N^{0.67}$ . Theories about phase separation in polymeric blends or in BCPs identify three distinct regimes:<sup>8</sup> the weak segregation limit (WSL), in which  $L_0 \propto N^{1/2}$ , the strong segregation limit (SSL), in which  $L_0 \propto N^{2/3}$ , and the super strong segregation limit (SSSL), in which  $L_0 \propto N^1$ . The scaling exponent obtained from the fit of the data reported in Figure 7a indicates that the PS-*b*-PMMA system is in the SSL regime in the full range of the investigated molar masses, thus confirming and extending the results obtained by Xu et al.<sup>24</sup> for  $M_n > 73 \text{ kg}\cdot\text{mol}^{-1}$ .

From a technological point of view, the fine control of the characteristic dimensions ( $d$ ,  $L_0$ ) of the cylinder morphology provides the possibility to fabricate a periodic nanoporous pattern with tunable pore density. On the basis of the measured  $L_0$  values we calculated the average pore density in the polymeric template for each BCP, reported in Figure 7b as a function of  $M_n$ . The calculated values range from  $0.2 \text{ Td}\cdot\text{inch}^{-2}$  for sample B205 to  $1.3 \text{ Td}\cdot\text{inch}^{-2}$  for sample B39, overcoming the  $1 \text{ Td}\cdot\text{inch}^{-2}$  barrier and providing a viable route for nanostructure patterning in applications requiring very high density.<sup>47</sup>

The intensity profiles extracted from the GISAXS patterns and reported in Figure 6a exhibit significant differences in the number of diffraction peaks and in their width, with the samples of intermediate  $M_n$  showing narrower and, thus, better resolved diffraction peaks. Indeed, the width of the peaks provides a quantitative measure of the in-plane structural order. Within the Hosemann's paracrystalline model,<sup>48</sup> the order correlation length  $\xi$  can be calculated as



**Figure 6.** (a) GISAXS intensity profiles for samples with different  $M_n$  measured along  $q_{xy}$  at fixed  $q_z$  ( $q_z = 0.22 \text{ nm}^{-1}$  for sample B39 and  $q_z = 0.24 \text{ nm}^{-1}$  for all other samples). These data were extracted from the GISAXS patterns reported in Figure 4. (b) GISAXS intensity profiles of sample B54 measured at three different incidence angles,  $\alpha_i = 0.12^\circ$ ,  $\alpha_i = 0.15^\circ$ , and  $\alpha_i = 0.18^\circ$ .

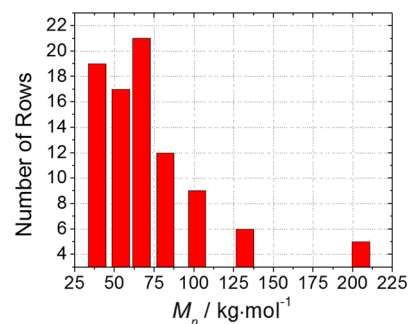


**Figure 7.** (a) Lattice spacing  $L_0$  measured by SEM (black circles) and GISAXS (red squares) as a function of the number of segments  $N$  in the BCP. (b) Pore density as a function of the  $M_n$ .

$$\xi = \frac{2\pi^3}{6.25 \cdot \Delta q_{10}}$$

where  $\Delta q_{10}$  is the integral half-width of the first order (10) diffraction peak. This quantity gives a measure of the distance over which the long-range positional order is statistically lost. Applied to our data, this model provides the  $\xi$  values reported in Table 2: the highest correlation length ( $\sim 700 \text{ nm}$ ) is obtained for sample B67, while the lowest ( $\sim 300 \text{ nm}$ ) is for sample B132 and sample B205.

To better understand the meaning of these results, it can be useful to consider the trend of the  $\xi$  values expressed in units of the lattice spacing, i.e., through the fraction  $\xi/L_0$  as a function of  $M_n$  (Figure 8). The histogram shows that for low  $M_n$  samples the in-plane order statistically extend over 18–20 consecutive lines of cylinders. With increasing  $M_n$  the fraction  $\xi/L_0$  decreases, reaching the minimum value of five lattice spacings for sample B205. Such result suggests that for high  $M_n$  values the 900 s annealing time is not sufficient to obtain a good ordering of the nanodomains for the direct application in standard photolithographic processes. Nevertheless, to increase the lateral ordering of the nanostructures and the registration with the underlying substrate, methods such as chemically patterned substrates<sup>49</sup> and graphoepitaxy<sup>50–52</sup> have been widely



**Figure 8.** Histogram reporting the ratio  $\xi/L_0$ , expressed as number of rows, as a function of the BCP  $M_n$ .

explored in the literature. In the latter case, the nanodomains are laterally confined within lithographically defined trenches whose widths are commensurate with  $L_0$ , forming perfectly organized polymeric structures that propagate over long ranges. The extension of nanodomains obtained on flat substrates well correlate with the typical trench width (normally between 50 and 500 nm) exploited in the literature.<sup>53,54</sup>

## CONCLUSIONS

In this work, the self-organization in PS-*b*-PMMA thin films undergoing a RTP treatment was investigated over a broad range of  $M_n$  between 29.4 kg·mol<sup>-1</sup> and 205 kg·mol<sup>-1</sup> as a function of the RTP annealing temperature. By optimizing this parameter for each  $M_n$ , we could get regular arrangements of hexagonally packed PMMA cylinders, embedded in a PS matrix and normally oriented with respect to the substrate, for all samples, except the B29 BCP. The hexagonally packed cylinder morphology, with  $d$  ranging between 12 and 30 nm and  $L_0$  between 24 and 73 nm, was obtained in 900 s keeping the thickness of the polymeric film constant at  $\sim$ 30 nm. The lateral ordering of the nanostructures, measured by GISAXS, extended over a maximum of 20 consecutive rows of PMMA cylinders for the BCP having low  $M_n$ . These results demonstrate the possibility to control the self-assembly process of BCPs over a wide range of  $M_n$  values with a very simple and flexible thermal treatment. It represents an important step toward the simplification of the processing conditions required for the integration of these materials in a standard photolithographic process flow.

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### Notes

The authors declare no competing financial interest.

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